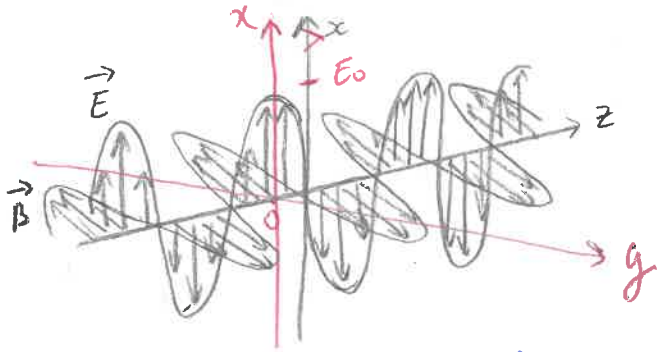


AD 1 TD 31, groupe 10.



à $t=0$ $z=0$ $\vec{E} = E_0 \vec{u}_x$ valeur max

Δ besoin de signe pour \vec{B} car à $t=0, z=0$
 $\vec{B} = \frac{E_0}{c} \vec{u}_y \rightarrow \text{max}$

$$\vec{E} = E_0 e^{j(\omega t - kz)} \vec{u}_x$$

→ dans le vide donc k réel.

$$\vec{B} = \frac{\vec{k} \wedge \vec{E}}{\omega} = \frac{k \vec{u}_z \wedge E_0 e^{j(\omega t - kz)} \vec{u}_x}{\omega} = \frac{k}{\omega} E_0 e^{j(\omega t - kz)} \vec{u}_y$$

Démonstration de $\frac{\vec{k} \wedge \vec{E}}{\omega} = \vec{B}$

$$\text{rot } \vec{E} = -\frac{\partial \vec{B}}{\partial t}$$

$$\vec{\nabla} \wedge \vec{E} = -\frac{\partial \vec{B}}{\partial t}$$

$$-jk \vec{u}_z \wedge \vec{E} = -j\omega \vec{B}$$

$$\frac{\vec{k} \wedge \vec{E}}{\omega} = \vec{B}$$

MF

coordonnées cartésiennes

OPPH

Δ relation non valable pour des ondes stationnaires.

revenir à $\vec{\nabla} \wedge \vec{E} = -\frac{\partial \vec{B}}{\partial t}$

avec $\vec{\nabla} \begin{pmatrix} \frac{\partial}{\partial x} \\ \frac{\partial}{\partial y} \\ \frac{\partial}{\partial z} \end{pmatrix}$

TD31 - AD2

groupe <

$$* I = \frac{P}{S} = \frac{3 \times 10^{-3} \text{ W}}{1 \times 10^{-6} \text{ m}^2} = 3 \times 10^3 \text{ W} \cdot \text{m}^{-2}$$

avec \vec{E}, \vec{B} OPPH

$$\text{on sait que } I = \langle \|\vec{\Pi}\| \rangle \text{ et } \vec{\Pi} = \frac{\vec{E} \wedge \vec{B}}{\mu_0}$$

$$* \|\vec{\Pi}\| = \left\| \frac{E_0 \cos(\omega t - kx) \vec{u}_y \wedge \frac{E_0}{c} \cos(\omega t - kx) \vec{u}_z}{\mu_0} \right\|$$

$$= \left\| \frac{E_0^2}{c \mu_0} \cos^2(\omega t - kx) \vec{u}_x \right\| = \frac{E_0^2}{c \mu_0} \cos^2(\omega t - kx)$$

$$* \langle \|\vec{\Pi}\| \rangle = \frac{E_0^2}{2c \mu_0}$$

$$* I = \langle \|\vec{\Pi}\| \rangle \Leftrightarrow E_0 = \sqrt{3 \times 10^3 \times 2 \times 3 \times 10^8 \times 4 \pi \times 10^{-7}} \\ = 1,5 \times 10^3 \text{ V} \cdot \text{m}^{-1}$$

$$\underline{E_0 = 1,5 \times 10^3 \text{ V} \cdot \text{m}^{-1}}$$

$$* B_0 = \frac{E_0}{c} = \frac{1,5 \times 10^3}{3 \times 10^8} = 5 \times 10^{-6} \text{ T}$$

$$\underline{B_0 = 5 \times 10^{-6} \text{ T}}$$

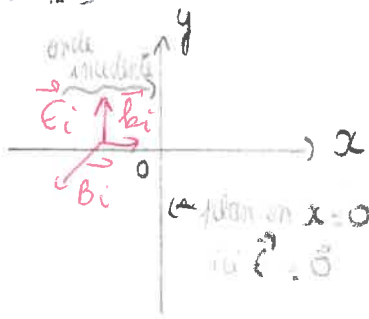
pas du cours

$$\left(* I = \langle u \rangle c \Leftrightarrow \langle u \rangle = \frac{I}{c} = \frac{P}{S \times c} = \frac{3 \times 10^{-3}}{10^{-6} \times 3 \times 10^8} = 10^{-5} \text{ J} \cdot \text{m}^{-3} \right)$$
$$\langle u \rangle = 10^{-5} \text{ J} \cdot \text{m}^{-3}$$

$$* N = \frac{P \times \lambda}{h c} = \frac{3 \times 10^{-3} \times 633 \times 10^{-9}}{3 \times 10^8 \times 6,62 \times 10^{-34}} \approx 9,6 \times 10^{15} \text{ photons}$$

$$* \langle u \rangle = \left\langle \frac{\epsilon_0 E^2}{2} + \frac{B^2}{2\mu_0} \right\rangle = \left\langle \frac{\epsilon_0 E_0^2 \cos^2(\omega t - kx)}{2} + \frac{B_0^2 \cos^2(\omega t - kx)}{2\mu_0} \right\rangle$$

$$= \frac{\epsilon_0 E_0^2}{4} + \frac{B_0^2}{4\mu_0} = \frac{1,5 E_0^2}{c^2 \mu_0 4} + \frac{E_0^2}{4 c^2 \mu_0} = \frac{E_0^2}{2 c^2 \mu_0} = \frac{(1,5 \times 10^3)^2}{2 (3 \times 10^8)^2 \times 4 \pi \times 10^{-7}} \approx 10^{-5} \text{ J} \cdot \text{m}^{-3}$$



$$\vec{E}_i = \epsilon_0 y e^{i(\omega t - kz + \phi_y)} \vec{u}_y$$

1) $\vec{E}_n = \epsilon_0 y e^{i(\omega t + kz + \phi_y)} \vec{u}_y$

En $x=0$: $\vec{E} = \vec{E}_i + \vec{E}_n = \vec{0} \Rightarrow \vec{E}_i = -\vec{E}_n$

$$\Leftrightarrow \epsilon_0 y e^{i(\omega t - 0 + \phi_y)} = -\epsilon_0 y e^{i(\omega t + 0 + \phi_y)}$$

$$\Leftrightarrow \epsilon_0 y e^{i\phi_y} = -\epsilon_0 y e^{i\phi_y} \quad \text{Vrai si } \epsilon_0 y = -\epsilon_0 y \text{ et } \phi_y = \phi_y$$

Donc $\vec{E} = \vec{E}_i + \vec{E}_n = \epsilon_0 y e^{i(\omega t - kz + \phi_y)} - \epsilon_0 y e^{i(\omega t + kz + \phi_y)}$

$$= \epsilon_0 y e^{i(\omega t + \phi_y)} (e^{-jkz} - e^{jkz}) \vec{u}_y$$

$$\vec{E} = -2\epsilon_0 y \sin(kz) e^{i(\omega t + \phi_y)} \vec{u}_y = 2\epsilon_0 y \sinh(kz) e^{i(\omega t + \phi_y - \pi/2)} \vec{u}_y$$

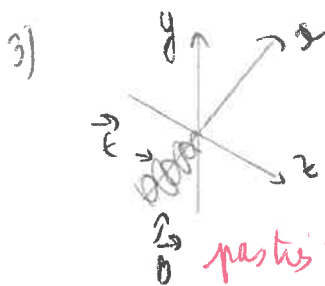
avec $-i = e^{-i\pi/2}$

2) Maxwell Faraday : $\text{rot } \vec{E} = -\frac{\partial \vec{B}}{\partial t}$

$$\vec{\nabla} \wedge \vec{E} = \begin{pmatrix} \frac{\partial}{\partial x} \\ \frac{\partial}{\partial y} \\ \frac{\partial}{\partial z} \end{pmatrix} \wedge \begin{pmatrix} 0 \\ 0 \\ E(z,t) \end{pmatrix} = \begin{pmatrix} -\frac{\partial}{\partial z} E(z,t) \\ 0 \\ \frac{\partial}{\partial x} E(z,t) \end{pmatrix}$$

Nous $-\frac{\partial \vec{B}}{\partial t} = -\frac{\partial(\mu_0 H)}{\partial t} \vec{u}_z + 0 + \frac{\partial(\epsilon_0 E)}{\partial t} \vec{u}_z$

$$= -2\epsilon_0 y j k \cos(kz) e^{i(\omega t + \phi_y)} \vec{u}_z \Rightarrow \vec{B} = \frac{+2\epsilon_0 y \cos(kz) e^{i(\omega t + \phi_y)}}{j\omega} \vec{u}_z = -\frac{2k\epsilon_0 y \cos(kz) e^{i(\omega t + \phi_y)}}{\omega} \vec{u}_z$$



Position des nœuds : x_N en $\vec{E} = \vec{0}$

$$\sin(kx_N) = 0 \Leftrightarrow kx_N = 0 + m\pi \Rightarrow x_{N_m} = \frac{m\pi}{k} = \frac{m\lambda}{2}$$

positions de \vec{B} x_{NB} $\cos(kx_{NB}) = 0$

$$kx_{NB} = \frac{\pi}{2} + p\pi \Rightarrow x_{NB_p} = \frac{\lambda}{4} + p\frac{\lambda}{2}$$

4) $\vec{E}(x=0^+) - \vec{E}(x=0^-) = \frac{\mathcal{T}}{\epsilon_0} \vec{u}_z$

$$\Rightarrow \vec{0} = \frac{\mathcal{T}}{\epsilon_0} \vec{u}_z \Rightarrow \mathcal{T} = 0$$

donc surface de charge

car elle fa en $x=0$.

$$\vec{B}(x=0^+) - \vec{B}(x=0^-) = \mu_0 \vec{j}_s \wedge \vec{u}_z = \mu_0 \int_s \vec{u}_y \wedge \vec{u}_z = \mu_0 \int_s \vec{u}_x = \frac{\mathcal{I}}{4} \vec{u}_x$$

$$\Rightarrow -\frac{2k\epsilon_0 y \cos(kz) e^{i(\omega t + \phi_y)}}{\omega} \vec{u}_z = \mu_0 \int_s \vec{u}_x \vec{u}_z$$

$$\Rightarrow \int_s = \frac{+2k\epsilon_0 y \cos(kz) e^{i(\omega t + \phi_y)}}{\mu_0 \omega} \vec{u}_y$$

onde se propageant dans
un milieu absorbant à la vitesse de phase

$$\begin{aligned} s(x, t) &= S_0 e^{-\frac{x}{\delta}} \cos(\omega t - kx) \\ &= S_0 e^{-\frac{x}{\delta}} \cos\left(\omega\left(t - \frac{k}{\omega}x\right)\right) \\ &= s\left(t - \frac{x}{v_p}\right) \end{aligned}$$

Par identification $v_p = \frac{\omega}{k}$ ✓

$$\begin{aligned} \underline{s}(x, t) &= \underline{S_0} e^{-\frac{x}{\delta}} e^{j(\omega t - kx)} \\ &= \underline{S_0} e^{-\frac{x}{\delta}} e^{j\left(\omega\left(t - \frac{x}{v_p}\right)\right)} \\ &= \underline{S_0} e^{j\left(\omega\left(t - \frac{x}{v_p}\right) - \frac{x}{\delta}\right)} \\ &= \underline{S_0} e^{j\left[\omega\left(t - \frac{x}{v_p}\right) + j\frac{x}{\delta}\right]} \\ &= \underline{S_0} e^{j\left[\omega t - x\left(\frac{\omega}{v_p} - j\frac{1}{\delta}\right)\right]} \end{aligned}$$

Par identification $\underline{k} = \frac{\omega}{v_p} - j\frac{1}{\delta}$

$$v_p \rightarrow s(x, t) = S_0 e^{-x/\delta} \cos\left(\omega\left(t - \frac{x}{v_p}\right)\right)$$

$$\begin{aligned} s(x, t) &= s\left(0, t - \frac{x}{v_p}\right) = s\left(t - \frac{x}{v_p}\right) \\ \text{car } v_p &= \frac{x}{\delta} \end{aligned}$$

$$\underline{S_0} e^{j(\omega t - kx)}$$

on cherche à exprimer
 \underline{k} en fonction de v_p et δ .